Amplitude analysis of $\gamma N \rightarrow K \pi \Lambda \& K p \rightarrow \pi \Lambda$

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Outline:

- Motivation
- Analysis of $\gamma N \rightarrow K \pi \Lambda$
- Analysis of $Kp \rightarrow \pi \Lambda$
- Summary

1. Motivation

Distinguishable model predictions for Σ^* of $1/2^-$ SU(3) octet					
	Quenched	&	unquenched quark models		
L=1	qqq excitation		L=0 qqqqq excitation		
udq	~ N*(1535)		udus s ~ N*(1535)		
uds	Λ*(1670), Σ*(~1630)	udsq \bar{q} $\Lambda^{*}(1405), \Sigma^{*}(\sim 1380)$		

Σ^* in PDG

****	$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$\Sigma^*(1670)3/2^-$ $\Sigma^*(2030)7/2^+$
***	$\Sigma^*(1660)1/2^+$ $\Sigma^*(1750)1/2^ \Sigma^*(2250)??$	Σ*(1940)3/2 ⁻
**	$\Sigma^{*}(1620)1/2^{-}$ $\Sigma^{*}(1690)??$ $\Sigma^{*}(2080)3/2^{+}$ $\Sigma^{*}(2455)??$ $\Sigma^{*}(2455)??$	E*(1880)1/2+ *(2620)??
*	$\begin{array}{llllllllllllllllllllllllllllllllllll$	*(1580)3/2 ⁻ ;*(2000)3/2 ⁻ *(3000)??

All from old experiments of 1970-1985 !!

No $\Sigma^*(1/2^-)$ around 1380 MeV ?

Re-analysis of old data on K⁻ p \rightarrow $\Lambda \pi^+ \pi^-$ Wu, Dulat, Zou, PRD80 (2009) 017503; PRC81 (2010) 045210

→ Possibly hidden $\Sigma^*(1/2^-)$ under $\Sigma^*(1385)3/2^+$ peak



New generation experiments on Σ^* at CLAS, LEPS, CB

CLAS 2005:
$$\gamma + p \to K^{+} + \Sigma^{*0}$$
 Ey=1.5-4 GeV
LEPS 2009: $\vec{\gamma} + n \to K^{+} + \Sigma^{*-}$ Ey=1.5-2.4 GeV
 $\downarrow \Lambda + \pi$
CB 2009: $K^{-} + p \to \pi^{0} + \Lambda$ Pk=514-750 MeV

Anything new on Σ^* ?

2. Analysis of $\gamma N \rightarrow K \pi \Lambda$

First studied with Effective Lagrangian approach by Oh, Ko, Nakayama, **PRC**77,045204(2008)



Feynman diagrams for $\gamma N \rightarrow K\Sigma^*(3/2^+)$

•Form factors and contact current:

$$F_M = \frac{\Lambda_M^2 - m_K^2}{\Lambda_M^2 - q_t^2},$$

Other channels

t-channel K exc.

$$F_B(q_{\rm ex}^2, M_{\rm ex}) = \left[\frac{n\Lambda_B^4}{n\Lambda_B^4 + (q_{\rm ex}^2 - M_{\rm ex}^2)^2}\right]^n$$

The contact current for $\gamma p \rightarrow K^+ \Sigma^{*0}$ is

•Haberzettl et. al, **PRC**74,045202(2006) $M_c^{\mu\nu} = ie \frac{f_{KN\Sigma^*}}{m_K} (g^{\mu\nu} f_t - q^{\mu} C^{\nu}),$ where $C^{\nu} = -(2q - k)^{\nu} \frac{f_t - 1}{t - m_K^2} [1 - h(1 - f_s)]$ $-(2p + k)^{\nu} \frac{f_s - 1}{s - M_N^2} [1 - h(1 - f_t)]$ The contact current for $\gamma n \rightarrow K^+ \Sigma^{*-}$ is

$$M_c^{\mu\nu} = ie\sqrt{2}\frac{f_{KN\Sigma^*}}{m_K}(g^{\mu\nu}f_t - q^{\mu}C^{\nu}),$$

where

$$f^{\nu} = -(2q-k)^{\nu} \frac{f_t - 1}{t - m_K^2} [1 - h(1 - f_u)] + (2p' - k)^{\nu} \frac{f_u - 1}{u - M_{\Sigma^*}^2} [1 - h(1 - f_t)]$$

where

$$f_t = F_M^2$$
 and $f_s = F_B^2(s, M_N)$
 $f_u = F_B^2(u, M_{\Sigma}^*)$

h is a free parameter to fit experiments.

Prediction vs data

- Total cross section $\gamma p \rightarrow K^+ \Sigma^{*0}$ of CLAS well described.
- differential cross section $\vec{\gamma}n \to K^+ \Sigma^{*-}$ of LEPS data also be described, but not for the Beam asymmetry A_{beam} .



Two possible solutions for the problem:





ds/dcos $\theta_{c.m.}$ for $\gamma n \rightarrow K\Sigma^*$ compared with LEPS data

Integrated cross section for $\gamma n \rightarrow K\Sigma^*$ vs. LEPS data

different predictions for the two schemes.

Scheme I describes both with same parameters, Scheme II shoud use different h . $\widehat{\underline{\mathfrak{A}}}$

Integrated cross section for $\gamma p \rightarrow K\Sigma^*$ vs. CLAS data



3. Analysis of Kp $\rightarrow \pi \Lambda$



The high precision new data can give valuable information for Σ^* resonances.





With these basic ingredients of 14 tunable parameters , the best fit gives $\chi^2 = 763$ for the 248 data points, including

Differential cross sections:
$$\frac{d\sigma_{\pi^{0}\Lambda}}{d\Omega} = \frac{d\sigma_{\pi^{0}\Lambda}}{2\pi d\cos\theta} = \frac{1}{64\pi^{2}s} \frac{|\mathbf{q}|}{|\mathbf{k}|} |\bar{\mathcal{M}}|^{2}$$
$$\Lambda \text{ Polarization :} \qquad P_{\Lambda} = \frac{3}{\alpha_{\Lambda}} \Big(\int \cos\theta' \frac{d\sigma_{K^{-}p \to \pi^{0}\Lambda \to \pi^{0}\pi N}}{d\Omega d\Omega'} d\Omega' \Big) \Big/ \frac{d\sigma_{\pi^{0}\Lambda}}{d\Omega}$$

Adding *** $\Sigma(1660)1/2^+$, $\chi^2 = 223$ for 248 data points with 18 tunable parameters.





$$K^- + p \rightarrow \pi^0 + \Lambda$$
 $\sqrt{s} = 1569 - 1676 \text{ MeV}$

Replacing $\Sigma(1635) 1/2^+$ by a $\Sigma(1/2^-)$, χ^2 increases by more than 160 with mass goes down to be below 1400 MeV.

 $\Sigma(1380)1/2^{-}$ is not needed, but cannot be excluded.

CB A Polarization data is crucial for discriminating $\Sigma^*(1620)1/2^-$ from $\Sigma(1635) 1/2^+$.

Summary

With the analysis of three reactions:

 $\gamma n \rightarrow K \Sigma^*$

K p→ ππΛ

The evidence of $\Sigma^*(1/2)$ predicted by the pentaquark models.

К р→πΛ

Existence of Σ(1660)1/2 +(***), with mass near 1635MeV, width 121MeV.

Need More experiment data to confirm them!!